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# HIGH-FREQUENCY ELECTROMAGNETIC WAVES DESTABILIZED BY RUNAWAY ELECTRONS IN A NEAR-CRITICAL ELECTRIC FIELD

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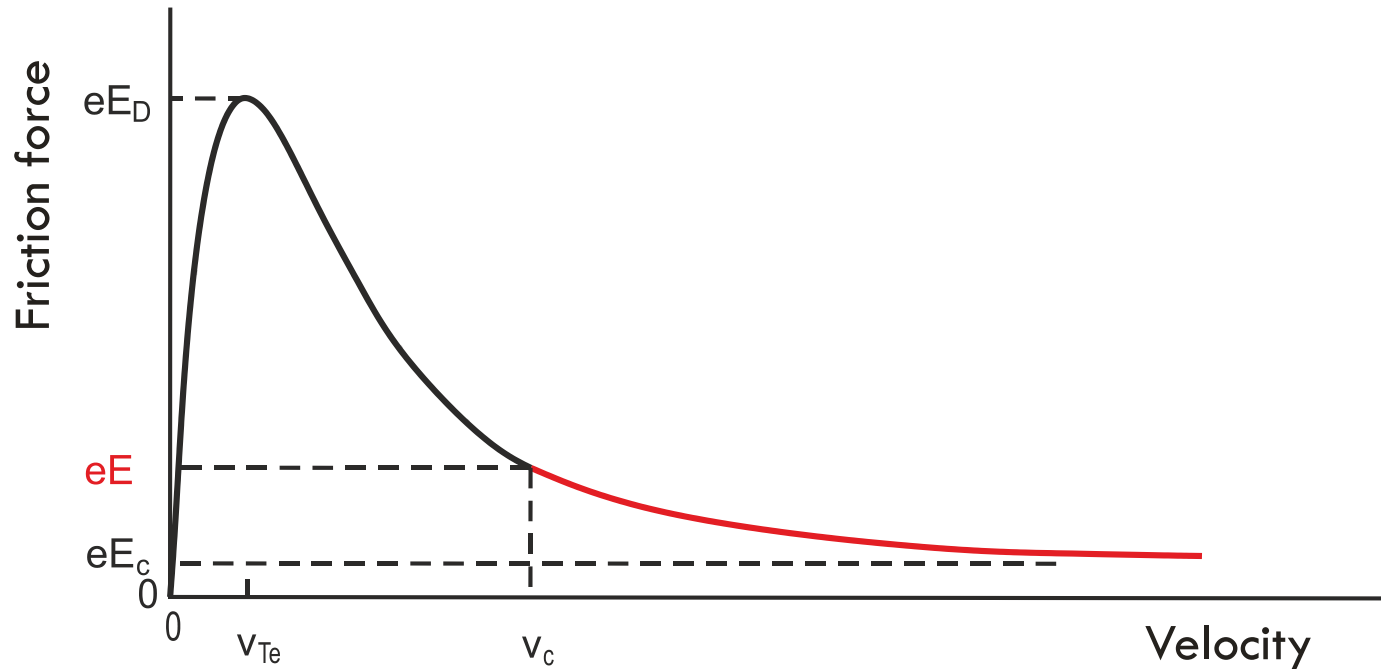
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# Runaway electrons

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- For electric fields higher than the critical field: accelerating force exceeds the friction force → *runaway electrons*

# Particle-wave interaction

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Dispersion relation of the plasma waves:

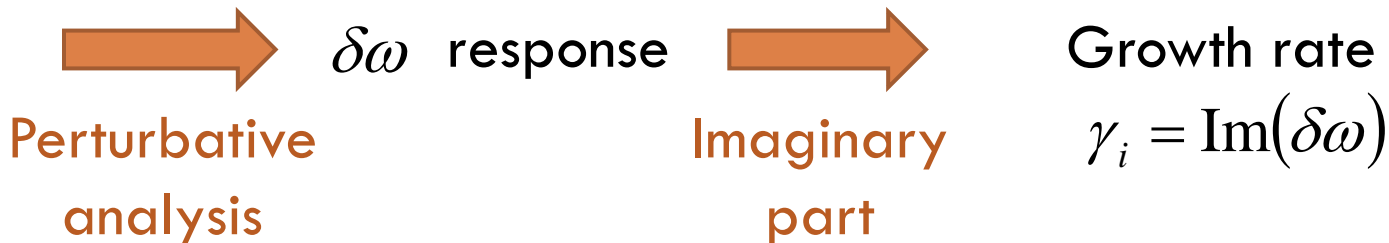
$$\left( \epsilon_{11} - \frac{k_{\parallel}^2 c^2}{\omega^2} \right) \cdot \left( \epsilon_{22} - \frac{k^2 c^2}{\omega^2} \right) + \epsilon_{12}^2 = 0$$

$$= \chi_{11}^r \left( \frac{k^2 c^2}{\omega^2} - \epsilon_{22} \right) + \chi_{22}^i \left( \frac{k_{\parallel}^2 c^2}{\omega^2} - \epsilon_{11} \right) + 2 \chi_{12}^r \cdot \epsilon_{12}$$

$\underline{\underline{\epsilon}} = \underline{\underline{1}} + \underline{\underline{\chi}}^e + \underline{\underline{\chi}}^i$

**Dielectric tensor**

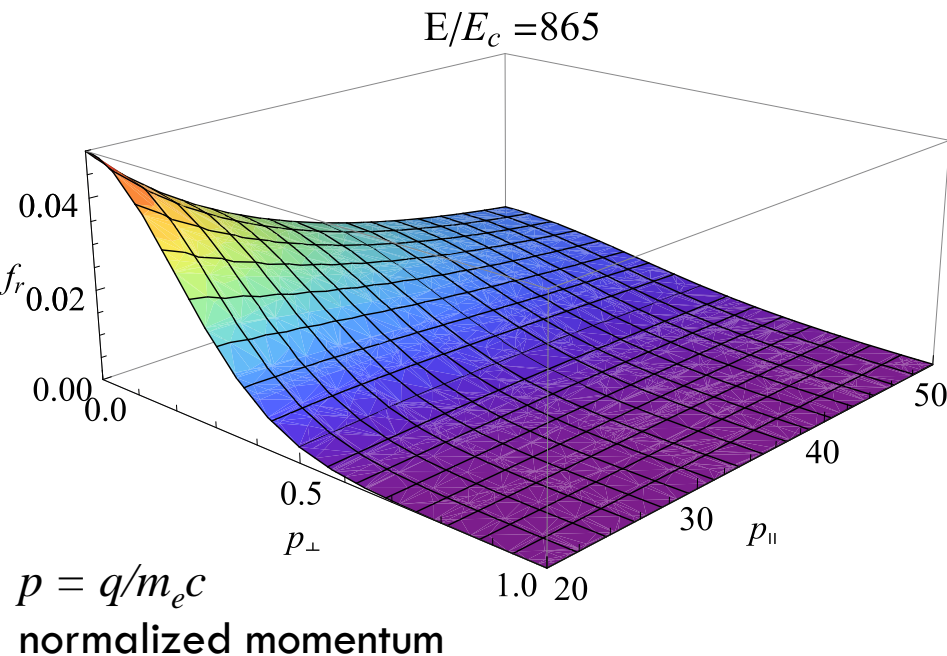
$k$ : wave number  
 $\omega$ : wave frequency  
 $c$ : speed of light



# Distribution function

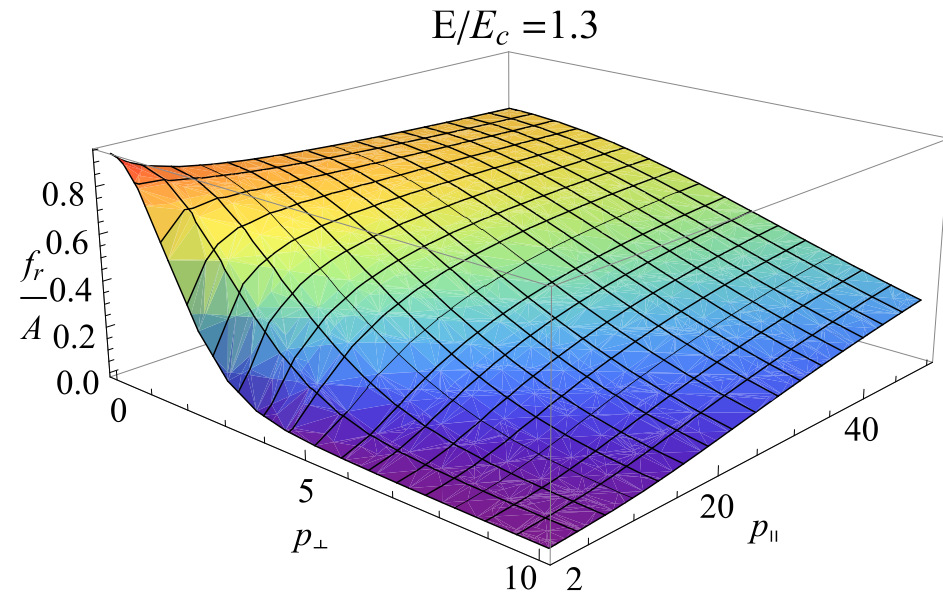
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High electric field



T. Fülöp, PoP **13**(062506), 2006

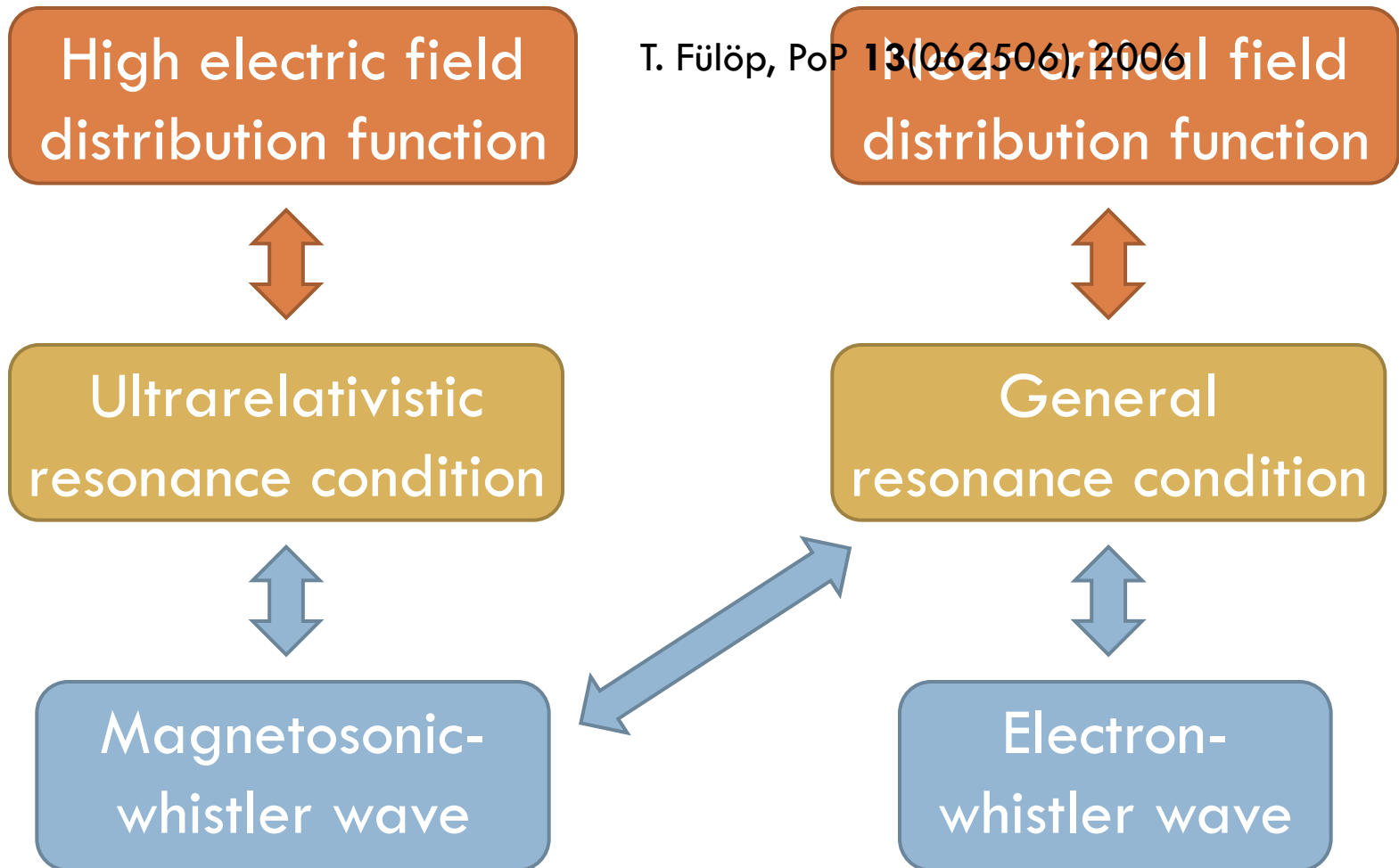
Near-critical field



P. Sandquist, PoP **13**(072108), 2006

# New approximations

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# Whistler waves

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## □ Electron-whistler wave

A. Kómár, BSc thesis, BME TTK, 2011

$$\omega_{ci} \ll \omega \quad \omega_{ce} \sqrt{m_e / m_i} \ll \omega$$

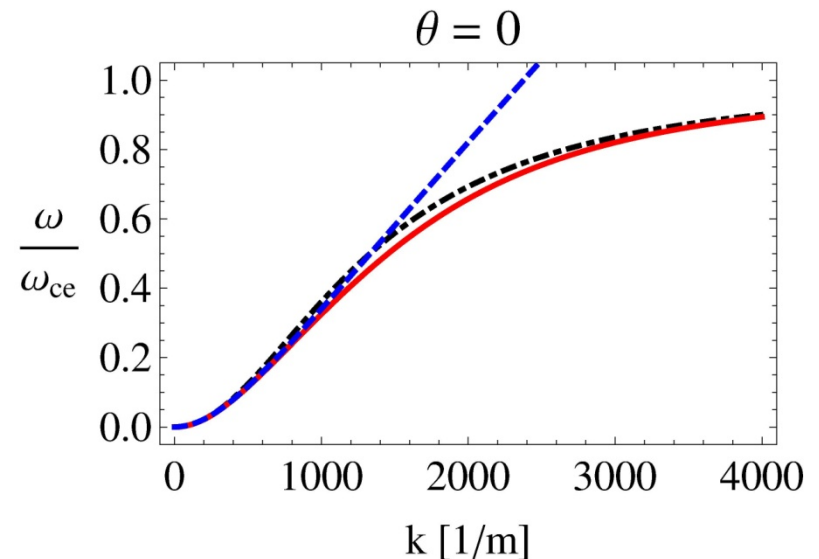
## □ Magnetosonic-whistler wave

T. Fülöp et al., PoP **13**(062506), 2006

$$\omega_{ci} \ll \omega \ll \omega_{ce}$$

## □ Whistler wave

S. Sazhin, Whistler-mode waves in a hot plasma, Cambridge Univ. Press, 1993

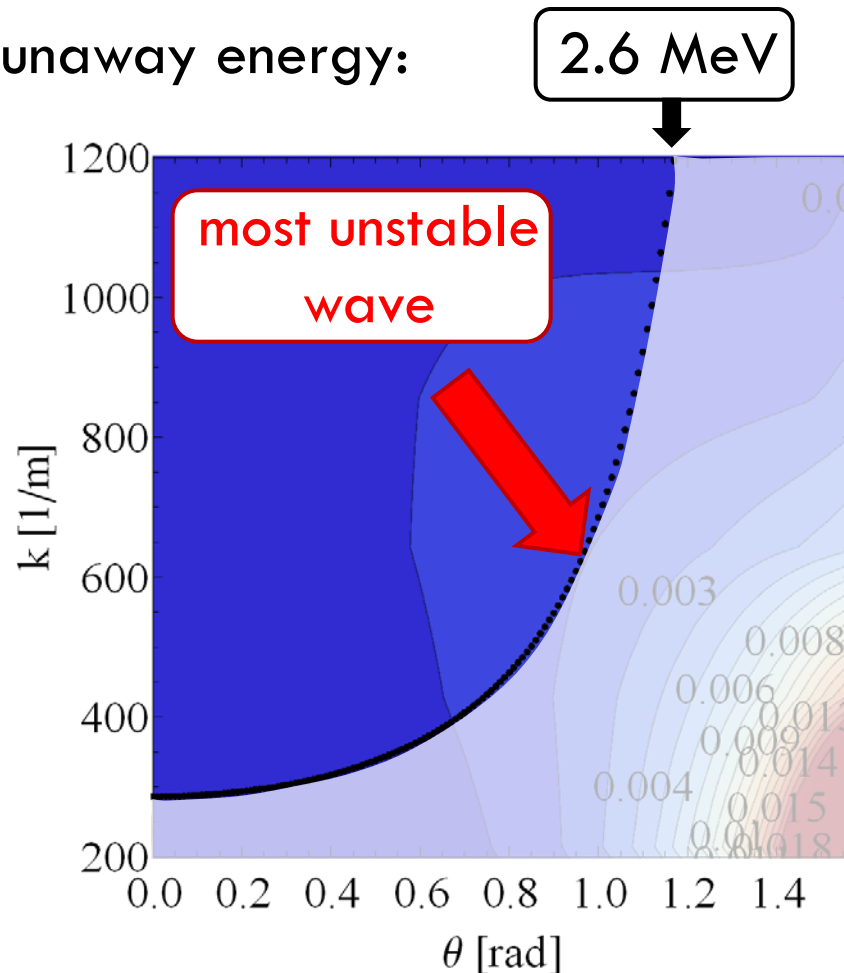




# Most unstable wave

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- Maximum runaway energy:



$$E/E_c = 1.3$$

$$B = 2 \text{ T}$$

$$n_e = 5 \cdot 10^{19} \text{ m}^{-3}$$

$$n_r = 3 \cdot 10^{17} \text{ m}^{-3}$$



# Damping rates of the wave, stability

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## □ Damping rates:

- **Collisional** damping: of electron-ion collisions
- **Convective** damping: the runaway beam has a finite radius,  $L_r$

## □ Stability:

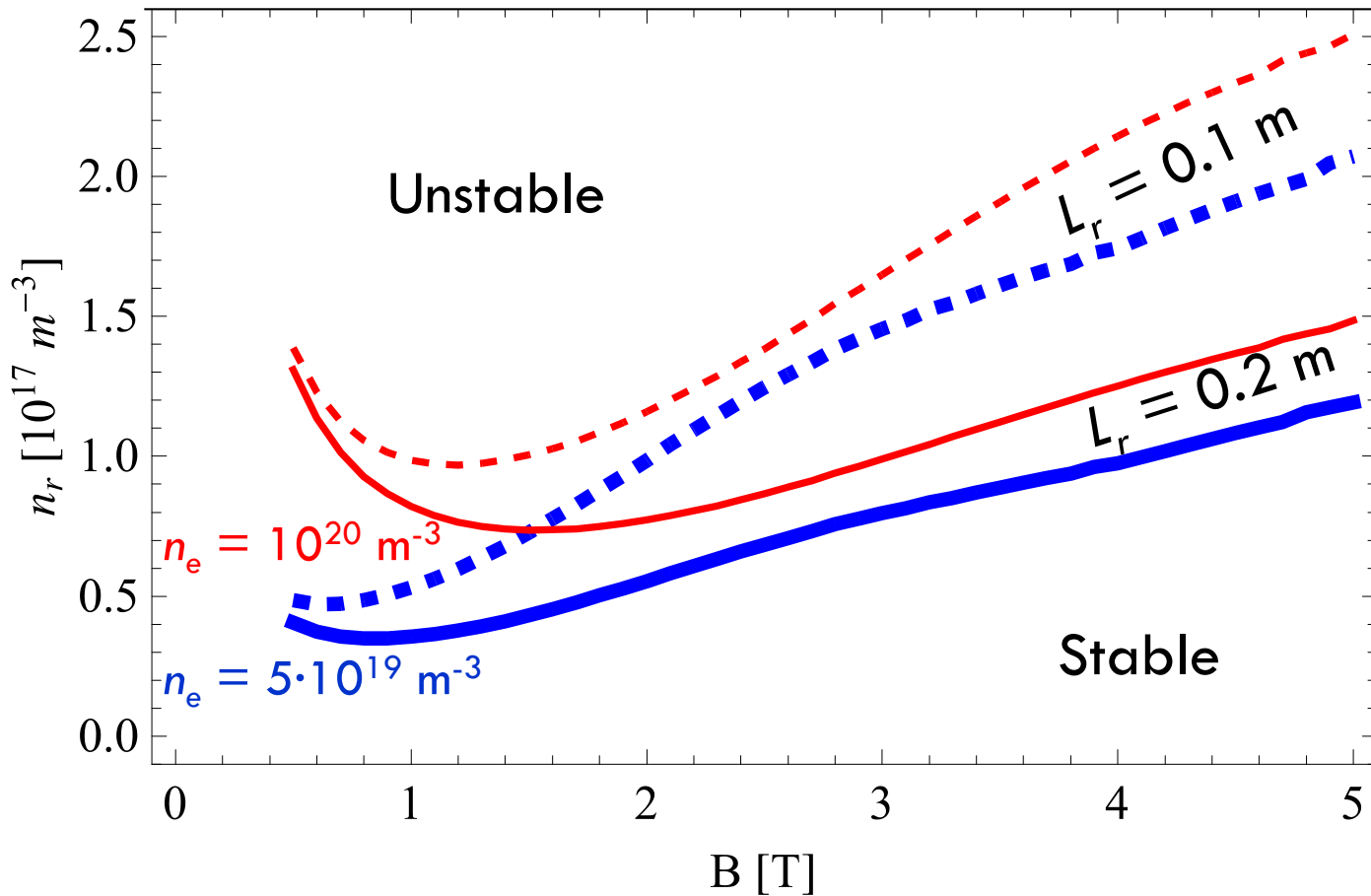
Finding (**Growth rate – Damping rates**) = 0  
(for the most unstable wave)



Critical runaway density

# Stability limit in near-critical field

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# Conclusions

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- Runaway electron – wave interaction in near-critical electric field
- Extending the previous approximations
  - ▣ General resonance condition
  - ▣ New whistler approximations
- Linear stability
  - ▣ The **most unstable wave** is an electron-whistler wave, dependent on the maximum runaway energy
  - ▣ **Stability threshold**: higher for higher magnetic field

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TÁMOP-4.2.2/B-10/1-2010-0009.

# General resonance condition

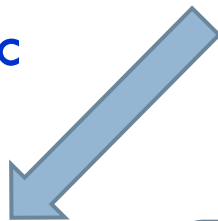
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- Runaway susceptibilities:  
calculated based on the distribution function,  $\iiint d^3p$
- **Implicit** resonance:

$$p_{par} + \frac{n\omega_{ce} - \omega_0 \cdot \gamma(p_{par}, p_{perp})}{k_{\parallel}c} = 0$$

Ultrarelativistic

$$\gamma \approx |p_{par}|$$



General case

$$\gamma = \sqrt{1 + p_{perp}^2 + p_{par}^2}$$



$$p_{res} = \frac{-n\omega_{ce}}{k_{\parallel}c - \omega_0}$$

$$p_{res} = \frac{-k_{\parallel}c n\omega_{ce} + \omega_0 \sqrt{(k_{\parallel}^2 c^2 - \omega_0^2)(1 + p_{perp}^2)} + n^2 \omega_{ce}^2}{k_{\parallel}^2 c^2 - \omega_0^2}$$

# Why the whistler wave?

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- Resonance condition is physical if:  $p_{res} > 0$
- In ultrarelativistic approximation:

$$p_{res} = \frac{-n\omega_{ce}}{k_{\parallel}c - \omega_0} \quad \Rightarrow \quad \boxed{k_{\parallel}c - \omega_0 > 0 \quad \text{and} \quad n < 0} \quad \text{or}$$
$$\boxed{k_{\parallel}c - \omega_0 < 0 \quad \text{and} \quad n > 0}$$

- In the general case:

$$\Rightarrow \quad \boxed{k_{\parallel}c - \omega_0(k, \theta) > 0} \quad \text{and} \quad n \leq 0$$



**Whistler wave**